



# Elastic modulus in the crystalline region and the thermal expansion coefficients of $\alpha$ -chitin determined using synchrotron radiated X-ray diffraction

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## ABSTRACT

Crystalline, orientated fibers of  $\alpha$ -chitin were obtained from the tendons of a snow crab (*Chionoecetes opilio*). The elastic modulus of the crystalline region along the fiber axis ( $E_l$ ) and the thermal expansion coefficients (TECs) of  $\alpha$ -chitin were measured using synchrotron X-ray diffraction. The  $E_l$  value of  $\alpha$ -chitin was  $59.3 \pm 11.3$  GPa, which was slightly larger than the value reported previously. This was because of the higher degree crystallinity of our sample and the use of synchrotron radiation as an X-ray source in this study. The linear TECs of the  $a$ -,  $b$ -, and  $c$ -axes were:  $\alpha_a = 6.1 \times 10^{-5} \text{ }^\circ\text{C}^{-1}$ ,  $\alpha_b = 7.3 \times 10^{-5} \text{ }^\circ\text{C}^{-1}$ , and  $\alpha_c = 2.1 \times 10^{-5} \text{ }^\circ\text{C}^{-1}$ , respectively. This anisotropic thermal expansion behavior,  $\alpha_b > \alpha_a > \alpha_c$ , was ascribed to the crystal structure and hydrogen bonding system of  $\alpha$ -chitin. In addition, the volume TEC of  $\alpha$ -chitin is reported for the first time as  $\beta = 15.8 \times 10^{-5} \text{ }^\circ\text{C}^{-1}$ .

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## 1. Introduction

Chitin, poly- $\beta(1 \rightarrow 4)$ -N-acetyl glucosamine, is synthesized by various organisms, such as insects, crustaceans, mollusks, fungi, yeasts, and algae (Rinaudo, 2006). Second only to cellulose, chitin is the next most abundant natural polymer on earth. Similar to cellulose, it exists in nature as crystalline fibrils that are a few nanometers across. There are two crystalline forms of natural chitin: the major form is  $\alpha$ -chitin, and the rarer form is  $\beta$ -chitin (Blackwell, 1973; Rudall, 1963). The crystal structure of these polymorphs has been investigated using X-ray diffraction, FT-IR spectroscopy, and solid-state  $^{13}\text{C}$  NMR spectroscopy (Blackwell, 1973; Focher, Naggi, Torri, Cosani, & Terbojevich, 1992; Rudall, 1963; Tanner, Chanzy, Vincendon, Roux, & Gaill, 1990). X-ray diffraction studies have established that these two polymorphs have different packing and arrangement of their molecular chains.  $\alpha$ -Chitin adopts a two-chain orthorhombic unit cell with  $P2_12_12_1$  symmetry, indicating an antiparallel chain arrangement (Minke & Blackwell, 1978; Sikorski, Hori, & Wada, 2009). On the other hand,  $\beta$ -chitin has a one-chain monoclinic unit cell, with  $P2_1$  symme-

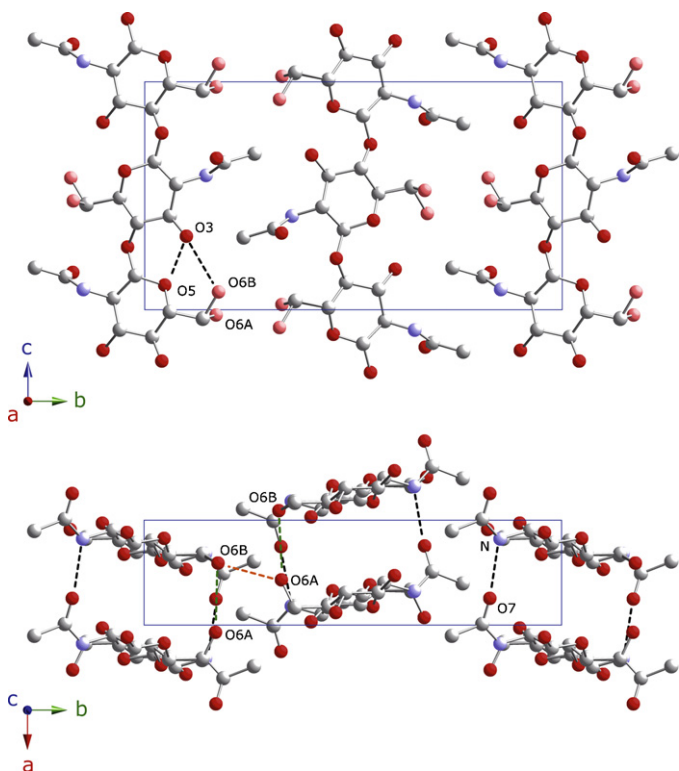
try, and consists of a parallel chain arrangement (Blackwell, 1969; Gardner & Blackwell, 1975).

It has been known that highly crystalline  $\alpha$ -chitin can be obtained from the spines of the phytoplankton, *Phaeocystis* (Chrétiennot-Dinet et al., 1997; Ogawa, Kimura, Wada, & Kuga, 2010) and the grasping spines of arrow worms (Atkins, Dlugosz, & Foord, 1979; Saito, Okano, Chanzy, & Sugiyama, 1995). Recently however, Sikorski et al. (2009) have reported on the relatively resolved and highly orientated X-ray fiber patterns from the tendons of a snow crab (*Chionoecetes opilio*), and thus, have reinvestigated the crystal structure of  $\alpha$ -chitin. The model they proposed includes two polymer chains with a  $2_1$  helix conformation, and an antiparallel arrangement in the orthorhombic unit cell ( $a = 0.475$  nm,  $b = 1.889$  nm, and  $c = 1.033$  nm) with a  $P2_12_12_1$  symmetry, the same as the classical model of Minke and Blackwell (1978). However, two distinctive conformations of the C6–O6 hydroxymethyl group, different from that suggested by Minke and Blackwell (1978), and the hydrogen bonding system of  $\alpha$ -chitin were demonstrated (Fig. 1).

The utilization of chitin is limited, despite its abundance in nature. To promote the use of chitin, it is important to investigate the structure/property relationships. However, little attention has been paid to this relationship for chitin. One of the important physical properties of the crystalline polymer is the elastic modulus of the crystalline region in the direction parallel to the chain axis ( $E_l$ ). The  $E_l$  value reflects the molecular conformation and the intramolecular hydrogen bonding system in the crystalline

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**Fig. 1.** Crystal structure and unit cell of  $\alpha$ -chitin viewed orthogonal to the  $bc$ -plane (top) and the  $ab$ -plane (bottom). The dotted lines denote hydrogen bonds (Sikorski et al., 2009).

lattice, and provides the maximum value for a sample modulus of the polymer. However, to the best of our knowledge, only a single paper has reported on the  $E_l$  value of  $\alpha$ -chitin (Nishino, Matsui, & Nakamae, 1999). The thermal expansion behavior is another important property that is related directly to the molecular arrangement and hydrogen bonding system in the crystalline lattice. However, the thermal expansion behavior reported for  $\alpha$ -chitin was only in the lateral directions (Wada & Saito, 2001).

In using synchrotron radiation as an X-ray source instead of conventional Cu K $\alpha$  radiation, the resolution of the data is enhanced with a much shorter exposure time, because a synchrotron can produce high intensity, directed, and highly monochromatized radiation. Therefore, in this study, the elastic modulus,  $E_l$ , and the thermal expansion behavior of  $\alpha$ -chitin crystals obtained from the tendons of a snow crab were investigated using synchrotron X-ray diffraction. The results obtained were compared with those obtained from another natural abundant polymer, cellulose.

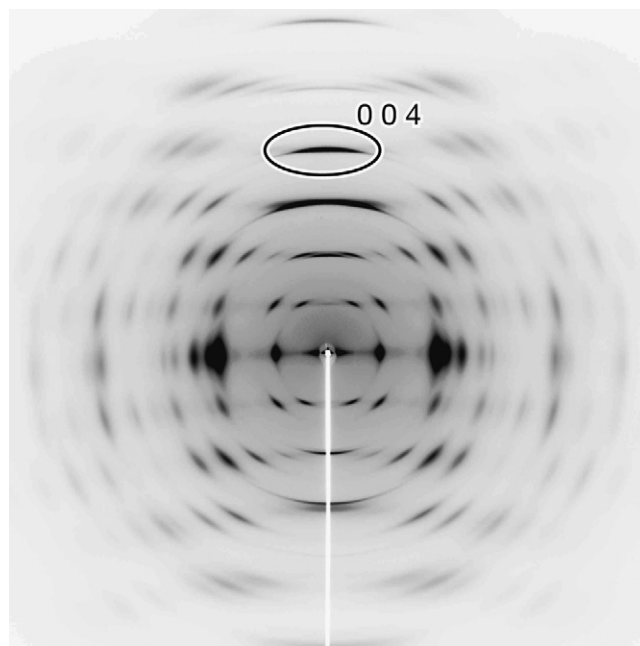
## 2. Experimental

### 2.1. Materials

The snow crab *C. opilio* was boiled in water for a period of 15 min, and the tendons were removed from its body. The tendons were purified by repetitive treatment with 5% KOH and 0.3% NaClO<sub>2</sub> solutions, as described previously (Sikorski et al., 2009). The samples were further boiled in a 0.1 N HCl solution for a period of 15 min, and washed in a large volume of deionized water. The purified tendons were finally dried in air while keeping the fiber length constant.

### 2.2. X-ray diffraction under tensile load

Both ends of the tendon fiber were fixed, with a span of 10 mm, to a pair of cardboard pieces using an epoxy adhesive. Eleven test



**Fig. 2.** Synchrotron X-ray fiber diffraction diagram obtained from the tendon of *C. opilio*. The longitudinal axis of the tendon was tilted to satisfy the Bragg condition for the meridional 004 diffraction.

samples were prepared and installed in the BL40B2 beam line at the Spring-8 synchrotron facility in Japan. The sample was clamped on a goniometer equipped with a stretching device and a load cell (LVS-2KA, Kyowa, Japan), and tilted from the fiber axis at an angle of 11.1° around an axis perpendicular to the fiber direction to comply with the Bragg condition for the meridional 004 diffraction. Synchrotron-radiated X-rays with a wavelength of 1.0 Å were irradiated on the sample for a period of 30 s at room temperature. A constant tensile load was applied to the sample during the measurements, and increased from 100 gf to 1900 gf in intervals of 100 gf (1 kgf = 9.8 N). The fiber diffraction patterns under set loads were recorded on imaging plates (IPs) (R-Axis IV, Rigaku, Japan). The sample-to-IP distance was calibrated using Si powder ( $d = 0.31355$  nm).

### 2.3. Degree of orientation

The degree of orientation of the microfibrils was evaluated using Hermans' order parameter (Alexander, 1979):

$$f = \frac{3\langle \cos^2 \phi \rangle - 1}{2} \quad (1)$$

where  $\langle \cos^2 \phi \rangle$  denotes the mean-square cosine of the angle between a given crystalline axis and the fiber axis. The parameter  $f$  was calculated from the azimuthal profile of the 004 reflection in the X-ray diffraction diagram of the tendon fiber with the tilted angle of 11.1° around an axis perpendicular to the fiber direction shown in Fig. 2, assuming a cylindrical symmetry.

### 2.4. Elastic modulus of the crystalline region

The lattice displacements were determined from the  $d$ -spacings of the 004 reflections ( $d_{004}$ ) from the recorded diffraction patterns. The meridional diffraction profiles were obtained using the Rigaku R-Axis software package. The peak-fitting routine of the X-ray diffraction profiles was carried out using a nonlinear least-squares fitting method (Wada, Okano, & Sugiyama, 1997). A pseudo-Voigt and linear functions were used to fit the 004 crystallite peaks and

the background profiles, respectively. The  $d_{004}$  value under a given load was calculated from the different peak positions. From the changes in the  $d_{004}$  value determined, the strain ( $\varepsilon$ ) was calculated using the following formula:

$$\varepsilon = \frac{\Delta d}{d_0} \quad (2)$$

where  $d_0$  is the  $d_{004}$  value at a load of 0 N for each sample, and  $\Delta d$  is the difference between the values of  $d_0$  and  $d_{004}$  under a given load.

The stress ( $\sigma$ ) in the crystalline region was assumed to be equal to the stress applied to the sample, as has been shown experimentally in previous work (Nishino et al., 1999; Nishino, Takano, & Nakamae, 1995). The cross-sectional area of the tendon fiber ( $A$ ) was measured as follows to evaluate the stress applied to the sample. The fiber was placed between polyethylene sheets and cut along its cross-sectional length. After coating the surface of the sample cross-section with  $\text{OsO}_4$ , images were recorded using an SEM (Hitachi S4000) and the area,  $A$ , was measured using an image analysis software package. The stress was then determined from the load divided by the value of  $A$  for each sample. Finally, the elastic modulus of the crystalline region ( $E_l$ ) was calculated as:

$$E_l = \frac{\sigma}{\varepsilon} \quad (3)$$

where  $\sigma$  is the stress applied to the sample, and  $\varepsilon$  is the strain.

## 2.5. X-ray diffraction at elevated temperatures

X-ray fiber diffraction at elevated temperatures was carried out at the BL40B2 beam line at the SPring-8 facility in Japan. The tendon fibers were mounted on a goniometer head with the incident X-ray beam orthogonal to the fiber axis. Synchrotron-radiated X-rays with a wavelength of 0.75 Å irradiated the sample under flowing nitrogen gas for a period of 60 s at a temperature between 50 and 250 °C in steps of 50 °C. The fiber patterns were recorded at the selected temperature using a camera system equipped with a flat imaging plate (IP) (R-Axis IV, Rigaku, Japan). The sample-to-IP distance was calibrated using Si powder ( $d = 0.31355$  nm).

From the recorded fiber patterns, the  $d$ -spacings of 28 diffractions were measured using the Rigaku R-Axis software package (Rigaku, Japan). These were indexed according to the orthogonal unit cell of Sikorski et al. (2009), and the unit-cell parameters were determined from the  $d$ -spacings and their indices using a least-squares fitting method. The linear and volume thermal expansion coefficients (TECs),  $\alpha$  and  $\beta$ , were determined using the following formulae:

$$\alpha = \frac{1}{l_{t=0}} \cdot \frac{\Delta l}{\Delta t} \quad (4)$$

$$\beta = \frac{1}{V_{t=0}} \cdot \frac{\Delta V}{\Delta t} \quad (5)$$

where  $l$  is either the  $d$ -spacing or the unit cell parameters,  $V$  is the volume of the unit cell, and  $t$  is the temperature (in °C).

## 3. Results and discussion

### 3.1. Elastic modulus

Fig. 2 shows the X-ray fiber diffraction diagram of a tendon of *C. opilio* at a tilt angle of 11.1° around the axis perpendicular to the fiber axis. The diagram in shown in Fig. 2 is a typical pattern of  $\alpha$ -chitin, and indicates the high crystallinity and high orientation of the microfibrils. The Hermans' ordered parameter,  $f_o$ , of the tendon was 0.89, evaluated from the meridional 004 reflection, which is almost the same value as ramie fiber in a naturally highly orientated

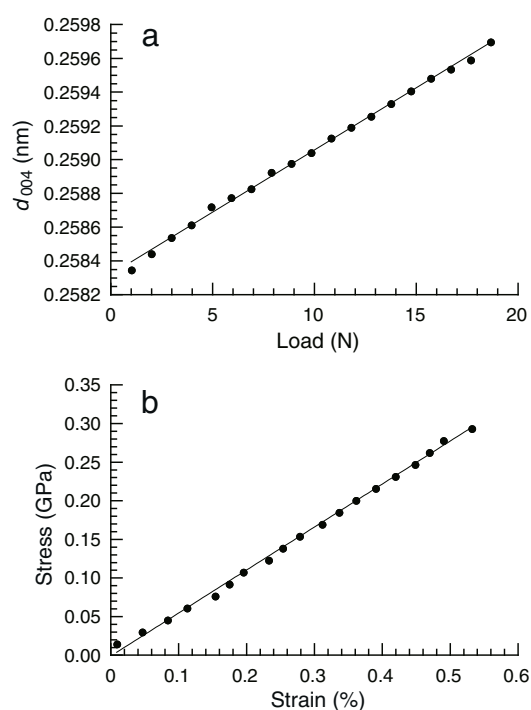
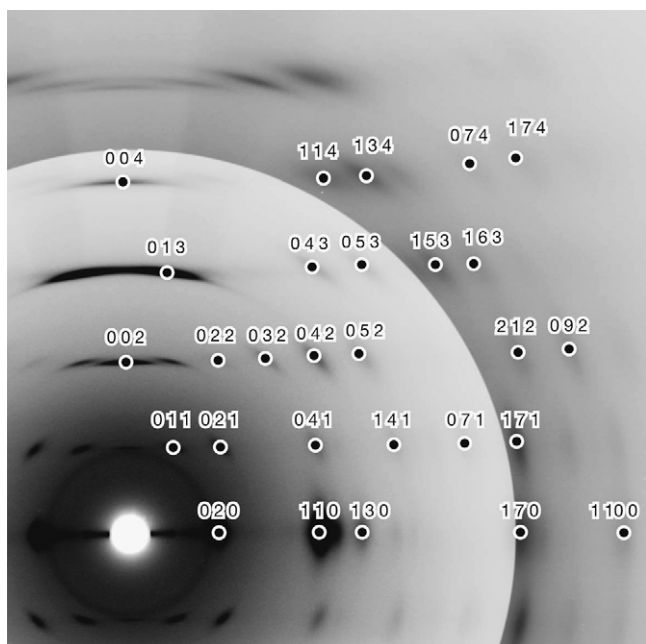


Fig. 3. (a) Changes in the  $d$ -spacings of the (004) plane,  $d_{004}$ , when applying a load to a typical tendon sample along its fiber axis. (b) The stress–strain curve obtained from the relationship between the  $d_{004}$  value and the load shown in part (a).

cellulose sample (Nishiyama, Kuga, Wada, & Okano, 1997). The high orientation of the microfibrils in our samples supported the internal stress homogeneity in the X-ray experiments under tensile load.

The change in  $d$ -spacing ( $d_{004}$ ) as a function of tensile load for a typical tendon sample is shown in Fig. 3a. The value of  $d_{004}$  increased with increasing load applied to the sample, and the relationship between  $d_{004}$  and the load was expressed as a straight line. This indicates that the experiment was conducted in the elastic deformation region. From the value of  $d_{004}$  and the load relationship (Fig. 3a), the stress–strain curve of the  $\alpha$ -chitin crystals shown in Fig. 3b was obtained. The stress of the  $\alpha$ -chitin crystals as a function of the tensile strain also showed a linear relationship over the experimental region, and the gradient of the line gave the elastic modulus,  $E_l$ , of the  $\alpha$ -chitin crystals. The average  $E_l$  value of the all the samples tested was  $59.3 \pm 11.3$  GPa, which is about 50% larger than that previously reported of 41 GPa (Nishino et al., 1999). This discrepancy is probably because the stress relaxation that interferes with the precise measurement of the lattice extension under load was restrained in using synchrotron X-ray radiation in this study. Moreover, the degree of crystallinity of the sample in this study was higher than that of the chitin fiber from a crab shell used in a previous paper (Nishino et al., 1999).

The  $E_l$  value of the  $\alpha$ -chitin in this study,  $59.3 \pm 11.3$  GPa, was less than half of the reported values of another naturally abundant polymer, cellulose, of 137–138 GPa (Nishino et al., 1995; Sakurada, Nukushina, & Ito, 1962). The cross-sectional area of a single molecule in the crystalline lattice ( $S$ ) of  $\alpha$ -chitin ( $44.9 \text{ Å}^2$ ) is larger than that of cellulose I $\beta$  ( $31.7 \text{ Å}^2$ ) (Nishiyama, Langan, & Chanzy, 2002; Sikorski et al., 2009). The large value of  $S$  of  $\alpha$ -chitin contributed to the small value of  $E_l$ , but the force needed to extend a single molecule of  $\alpha$ -chitin and cellulose I $\beta$  had a ratio of 1:1.6. This difference could be considered to be from the molecular conformation and intramolecular hydrogen bonding system of  $\alpha$ -chitin and cellulose I $\beta$  (Nishiyama et al., 2002; Sikorski et al., 2009). The fiber repeat of  $\alpha$ -chitin (10.33 Å) is smaller than that of cellulose I $\beta$  (10.38 Å), and therefore, the bond angle of the glucosidic linkage is



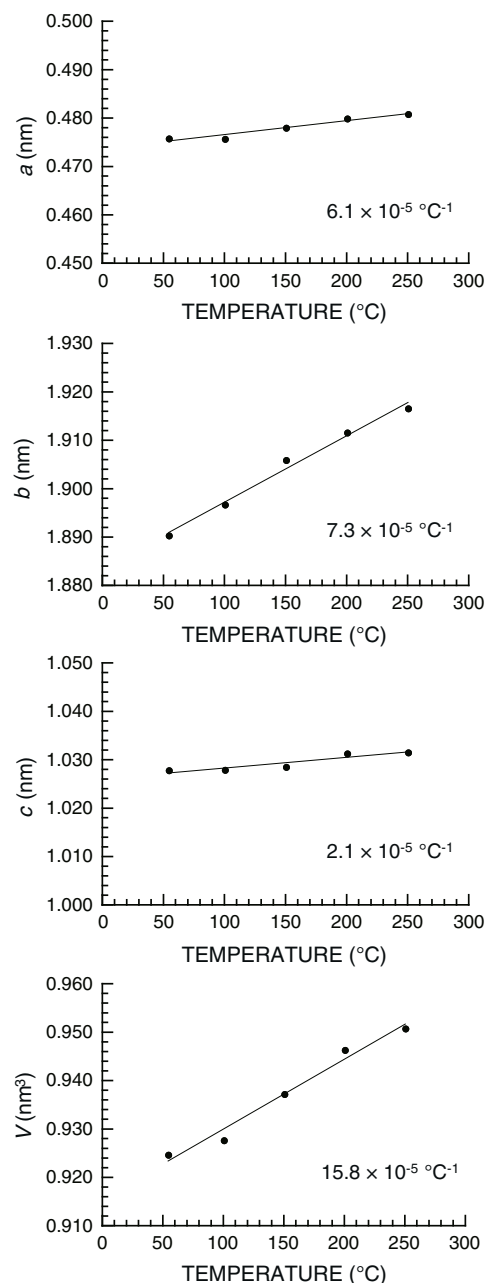
**Fig. 4.** Synchrotron X-ray fiber diffraction diagram of the tendon of *C. opilio* recorded at 50 °C. Twenty-nine reflections, together with indices used for the calculation of the unit cell parameters, are also shown. The image contrast of the outer region is enhanced to show the weaker reflections more clearly.

smaller: 113.96° for  $\alpha$ -chitin and 115.05° for cellulose I $\beta$ . Since the molecular chains of  $\alpha$ -chitin are more bent, the crystalline lattice easily expands along the fiber axis direction. There is only one bifurcated intramolecular hydrogen bond, O3–H $\cdots$ O5/O6B, in  $\alpha$ -chitin. On the other hand, in cellulose I $\beta$ , two hydrogen bonds, O3–H $\cdots$ O5 and O2–H $\cdots$ O6 or O6–H $\cdots$ O2, resist any extension along the fiber axis direction.

### 3.2. Thermal expansion

Synchrotron X-ray fiber diffraction diagrams of  $\alpha$ -chitin at given temperatures from 50 to 250 °C were recorded. The fiber pattern at 50 °C with markings for 28 diffractions and their indices used for analyses of the thermal expansion behavior is shown in Fig. 4. The unit cell parameters were calculated using a least-squares refinement of the  $d$ -spacings of these reflections without restraining the included angle of the  $a$ - and  $b$ -axes,  $\gamma$ , on taking into account the change in the angle with heating. However, the value of  $\gamma$  was almost constant at 90°; the change in the angle  $\gamma$  between 50 and 250 °C was within 0.2°. This result indicates that the unit cell is consistently orthorhombic in this temperature region. Therefore, the unit cell parameters,  $a$ ,  $b$ , and  $c$  at given temperatures, were determined based on the orthorhombic unit cell. These parameters and the volume of the unit cell ( $V$ ) as a function of the temperature are shown in Fig. 5.

The  $a$ -,  $b$ -, and  $c$ -axes increased gradually and linearly with increasing temperature. The linear TECs calculated using Eq. (4) were  $\alpha_a = 6.1 \times 10^{-5} \text{ }^\circ\text{C}^{-1}$ ,  $\alpha_b = 7.3 \times 10^{-5} \text{ }^\circ\text{C}^{-1}$ , and  $\alpha_c = 2.1 \times 10^{-5} \text{ }^\circ\text{C}^{-1}$ , respectively. This anisotropic thermal expansion behavior among the  $a$ -,  $b$ -, and  $c$ -axes,  $\alpha_b > \alpha_a > \alpha_c$ , is ascribed to the crystal structure and to the hydrogen bonding system of  $\alpha$ -chitin.  $\alpha$ -Chitin greatly expands along the  $b$ -axis direction, because there is only a single intermolecular hydrogen bond, O6A–H $\cdots$ O6B, which exists along the direction nearly parallel to this direction: i.e., the [110] direction. Meanwhile, the unit cell expands less along the  $a$ -axis direction, because two intermolecular hydrogen bonds, N–H $\cdots$ O7 and O6B–H $\cdots$ O6A, exist along this direction. The



**Fig. 5.** Changes in the unit cell parameters and volume of  $\alpha$ -chitin during heating from 50 to 250 °C under flowing nitrogen. The linear and volume TECs of the  $\alpha$ -chitin unit cells are also shown.

values of the TECs were the maximum at the  $b$ -axes direction and the minimum at the  $a$ -axes direction in lateral plane of  $\alpha$ -chitin crystal. The anisotropy of the TECs of  $\alpha$ -chitin is most notable between these two directions. However, this anisotropy in the lateral directions was not remarkable compared with cellulose:  $\alpha_b/\alpha_a = 1.2$  and  $\alpha_a/\alpha_b = 8.2$  for  $\alpha$ -chitin and cellulose I $\beta$ , respectively (Wada, 2002). In cellulose I $\beta$ , all the intermolecular hydrogen bonds exist along the  $b$ -axis direction, but there are no hydrogen bonds along the  $a$ -axis direction.  $\alpha$ -Chitin hardly expands along the  $c$ -axis direction, because the extended molecular chains run parallel to this direction, and a bifurcated intramolecular hydrogen bond, O3–H $\cdots$ O5/O6B, further stabilizes the molecule. However, the TEC value was 3.5 times larger than that of cellulose I $\beta$  (Wada, 2002). This is probably because of the fact that two intramolecular hydrogen bonds exist along the  $c$ -axis direction in cellulose I $\beta$ . In



addition, the fiber repeat of  $\alpha$ -chitin is slightly smaller than that of cellulose I $\beta$  (Nishiyama et al., 2002), and therefore, the molecular chains of  $\alpha$ -chitin may have room for extending on heating.

The unit-cell volume ( $V$ ) also showed a linear expansion with increasing temperature (Fig. 5). The volume TEC was calculated using Eq. (5) to be  $\beta = 15.8 \times 10^{-5} \text{ } ^\circ\text{C}^{-1}$ . This is approximately 1.4 times larger than the reported value of cellulose I $\beta$  of  $11.1 \times 10^{-5} \text{ } ^\circ\text{C}^{-1}$  (Hori & Wada, 2005). The volume occupied by a single molecule in the crystalline lattice is smaller, thus  $\alpha$ -chitin can expand more easily than can cellulose I $\beta$ .

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